

The Why,
What, Where
of the Long
Baseline
Neutrino
Experiment

Mary Bishai
Brookhaven
National
Laboratory

Why
Neutrinos:
History of ν s
Neutrino Mixing
Long Baseline ν
Oscillations

Accelerator ν
Experiments
Beams
Detectors

Which
Baseline?

Whats LBNE?

Summary and
Conclusions

The Why, What, Where of the Long Baseline Neutrino Experiment

INPA, Lawrence Berkeley National Lab, 03/01/2013

Mary Bishai
Brookhaven National Laboratory

March 1, 2013

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- 1 Why Neutrinos:
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- 2 Accelerator ν Experiments
 - Beams
 - Detectors
- 3 Which Baseline?
- 4 Whats LBNE?
- 5 Summary and Conclusions

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A BRIEF HISTORY OF NEUTRINOS

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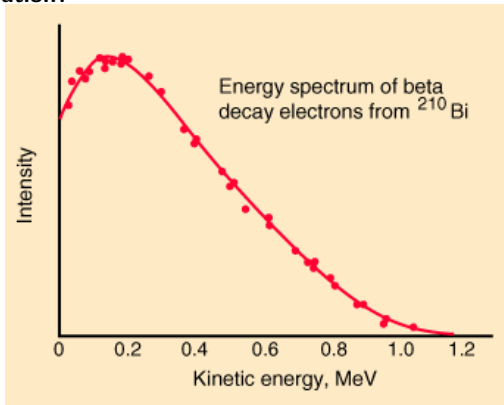
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Before 1930's: beta decay spectrum continuous - is this energy non-conservation?



G.J Neary, Proc Phys. Soc., A175, 71 (1940)

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Dec 1930: Wolfgang Pauli's letter to physicists at a workshop in Tübingen proposes that a neutrally charged "neutron" with a mass " < 0.01 proton mass" is emitted in beta-decays.

Dear Radioactive Ladies and Gentlemen,



Wolfgang Pauli

As the bearer of these lines, to whom I graciously ask you to listen, will explain to you in more detail, how because of the "wrong" statistics of the N and Li6 nuclei and the continuous beta spectrum, I have hit upon a desperate remedy to save the "exchange theorem" of statistics and the law of conservation of energy. Namely, the possibility that there could exist in the nuclei electrically neutral particles, that I wish to call neutrons, which have spin 1/2 and obey the exclusion principle and which further differ from light quanta in that they do not travel with the velocity of light. The mass of the neutrons should be of the same order of magnitude as the electron mass and in any event not larger than 0.01 proton masses. The continuous beta spectrum would then become understandable by the assumption that in beta decay a neutron is emitted in addition to the electron such that the sum of the energies of the neutron and the electron is constant...

I agree that my remedy could seem incredible because one should have seen those neutrons very earlier if they really exist. But only the one who dare can win and the difficult situation, due to the continuous structure of the beta spectrum, is lighted by a remark of my honoured predecessor, Mr Debye, who told me recently in Bruxelles: "Oh, It's well better not to think to this at all, like new taxes". From now on, every solution to the issue must be discussed. Thus, dear radioactive people, look and judge.

Unfortunately, **I cannot appear in Tübingen personally since I am indispensable here in Zurich because of a ball on the night of 6/7 December.** With my best regards to you, and also to Mr Back.

Your humble servant

. W. Pauli

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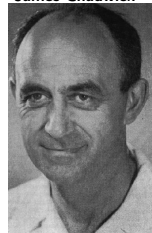
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1932: James Chadwick discovers the neutron -
its too heavy - cant be Pauli's particle



James Chadwick



Enrico Fermi

Solvay Conference, Bruxelles 1933: Fermi
proposes to name Pauli's particle the "neutrino".

The Theory of Weak Interactions

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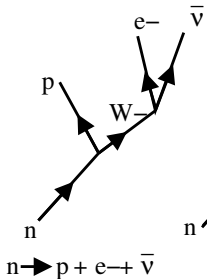
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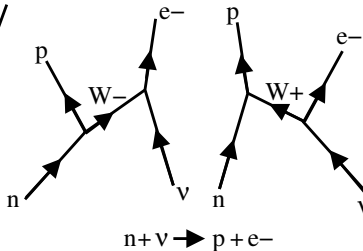
≥ 1933 : Development of the theory of **weak interactions and beta decay**

Charged current interactions

Decay of neutron

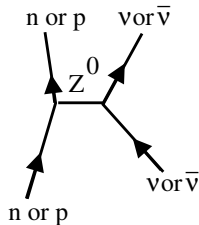


Neutrino interacts
with neutron



Neutral current interactions

n or p interacts with
neutrino or antineutrino



Finding Neutrinos...

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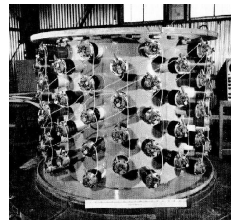
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1950's: **Fred Reines at Los Alamos and Clyde Cowan** mounted an experiment at the Hanford nuclear reactor in 1953 and in 1955 at the new Savannah River nuclear reactor. A detector filled with **water with CdCl_2 in solution** was located 11 meters from the reactor center and 12 meters underground.

The detection sequence was as follows:

- 1 $\bar{\nu}_e + p \rightarrow n + e^+$
- 2 $e^+ + e^- \rightarrow \gamma\gamma$ (2X 0.511 MeV + T_e^+)
- 3 $n + {}^{108}\text{Cd} \rightarrow {}^{109}\text{Cd}^* \rightarrow {}^{109}\text{Cd} + \gamma$
($\tau = 5\mu\text{s}$).



Neutrinos first detected using a nuclear reactor!

Neutrinos have Flavors

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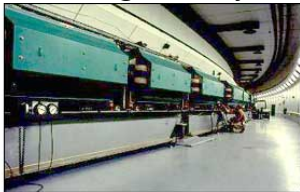
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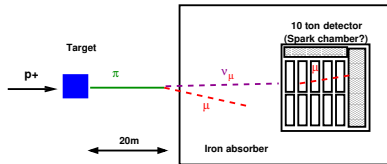
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1962: Leon Lederman, Melvin Schwartz and Jack Steinberger use BNL's Alternating Gradient Synchrotron (AGS) to produce a beam of neutrinos using the decay $\pi \rightarrow \mu \nu_x$



The AGS



Making ν 's

Result: 40 neutrino interactions recorded in the detector, 6 of the resultant particles were identified as background and 34 identified as

$$\mu \Rightarrow \nu_x = \nu_\mu$$

The first accelerator neutrino experiment was at the AGS.

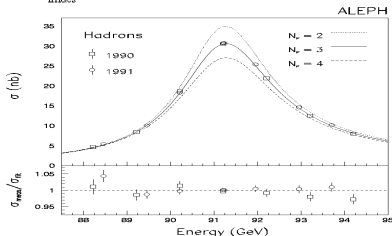
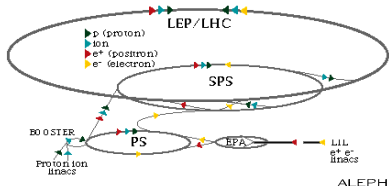
Number of Neutrino Flavors: Particle Colliders

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1980's - 90's: The number of neutrino types is precisely determined from studies of Z^0 boson properties produced in e^+e^- particle colliders. $N_\nu = 2.984 \pm 0.008$

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The LEP e^+e^- collider at CERN, Switzerland



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Direct Observation of ν_τ

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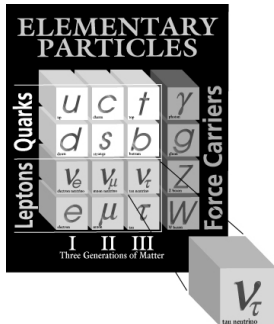
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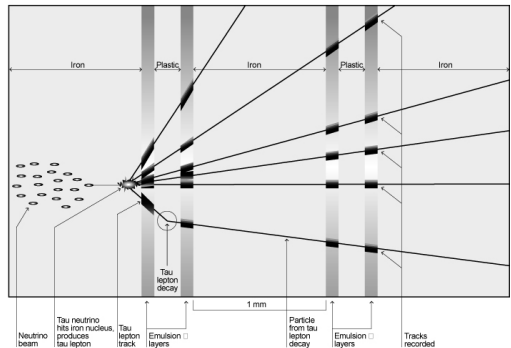
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July 20, 2000. The DONUT experiment finds evidence for the 3rd neutrino:



Detecting a Tau Neutrino



Of one million million tau neutrinos crossing the DONUT detector, scientists expect about one to interact with an iron nucleus.

Standard model: 3rd neutrino is the partner of the τ lepton

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NEUTRINO MIXING AND OSCILLATIONS

Neutrino oscillations

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1957,1967: **B. Pontecorvo** proposes that neutrinos could oscillate:

$$\begin{pmatrix} \nu_a \\ \nu_b \end{pmatrix} = \begin{pmatrix} \cos(\theta) & \sin(\theta) \\ -\sin(\theta) & \cos(\theta) \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

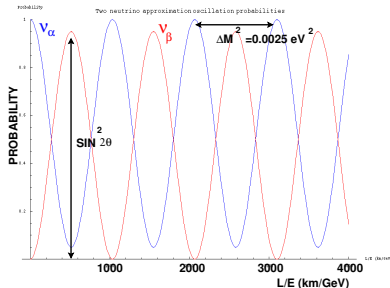
$$\begin{aligned} \nu_a(t) &= \cos(\theta)\nu_1(t) + \sin(\theta)\nu_2(t) \\ P(\nu_a \rightarrow \nu_b) &= |\langle \nu_b | \nu_a(t) \rangle|^2 \\ &= \sin^2(\theta) \cos^2(\theta) |e^{-iE_2 t} - e^{-iE_1 t}|^2 \end{aligned}$$

$$P(\nu_a \rightarrow \nu_b) = \sin^2 2\theta \sin^2 \frac{1.27 \Delta m_{21}^2 L}{E}$$

where $\Delta m_{21}^2 = (m_2^2 - m_1^2)$ in eV^2 ,
 L (km) and E (GeV).

Observation of oscillations

implies non-zero mass eigenstates



The Homestake Experiment

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1967: **Ray Davis** from BNL installs a large detector, containing 615 tons of tetrachloroethylene (cleaning fluid), 1.6km underground in Homestake mine, SD.

1 $\nu_e^{\text{sun}} + {}^{37}\text{Cl} \rightarrow e^- + {}^{37}\text{Ar}$, $\tau({}^{37}\text{Ar}) = 35$ days.

2 Number of Ar atoms \approx number of ν_e^{sun} interactions.



Ray Davis

Results: 1969 - 1993 Measured 2.5 ± 0.2 SNU (1 SNU = 1 neutrino interaction per second for 10^{36} target atoms) while theory predicts 8 SNU. This is a **ν_e^{sun} deficit of 69%**.

Solar ν_e disappearance \Rightarrow

first experimental hint of oscillations

SNO Experiment: Solar $\nu_e \rightarrow \nu_x$ Measurements

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2001-02: Sudbury Neutrino Observatory. Water Čerenkov detector with 1 kT heavy water (**0.5 B\$ worth on loan from Atomic Energy of Canada Ltd.**) located 2Km below ground in INCO's Creighton nickel mine near Sudbury, Ontario. Can detect the following ν^{sun} interactions:

- 1) $\nu_e + d \rightarrow e^- + p + p$ (CC).
- 2) $\nu_x + d \rightarrow p + n + \nu_x$ (NC).
- 3) $\nu_x + e^- \rightarrow e^- + \nu_x$ (ES).

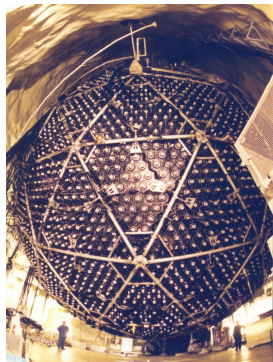
SNO measured:

$$\phi_{\text{SNO}}^{\text{CC}}(\nu_e) = 1.75 \pm 0.07(\text{stat})_{-0.11}^{+0.12}(\text{sys.}) \pm 0.05(\text{theor}) \times 10^6 \text{cm}^{-2} \text{s}^{-1}$$

$$\phi_{\text{SNO}}^{\text{ES}}(\nu_x) = 2.39 \pm 0.34(\text{stat})_{-0.14}^{+0.16}(\text{sys.}) \pm \times 10^6 \text{cm}^{-2} \text{s}^{-1}$$

$$\phi_{\text{SNO}}^{\text{NC}}(\nu_x) = 5.09 \pm 0.44(\text{stat})_{-0.43}^{+0.46}(\text{sys.}) \pm \times 10^6 \text{cm}^{-2} \text{s}^{-1}$$

All the solar ν 's are there but ν_e appears as ν_x !



The Super-Kamiokande Experiment. Kamioka Mine, Japan

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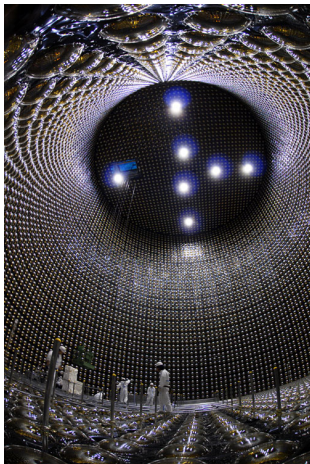
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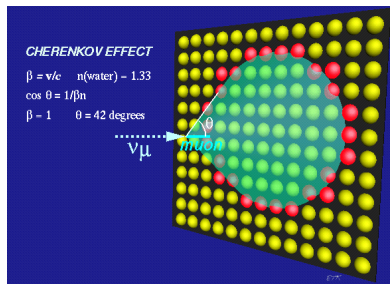
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50kT double layered tank of ultra pure water surrounded by 11,146 20" diameter photomultiplier tubes.

Neutrinos are identified by using CC interaction $\nu_{\mu,e} \rightarrow e^{\pm}, \mu^{\pm} X$. The lepton produces Cherenkov light as it goes through the detector:



The Super-Kamiokande Experiment. Kamioka Mine, Japan

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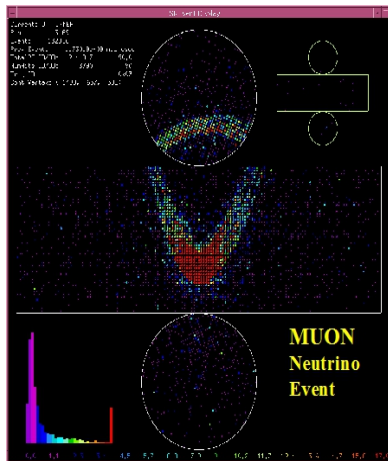
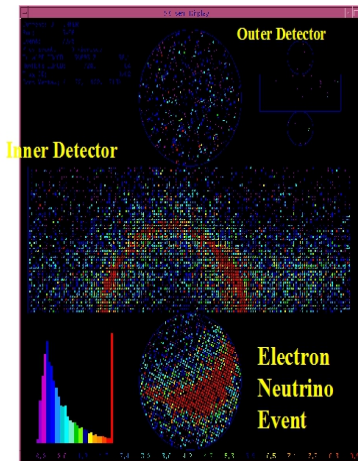
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Atmospheric Neutrino Oscillations: $\nu_\mu, \nu_e, \bar{\nu}$

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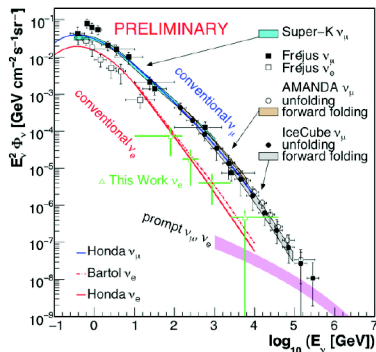
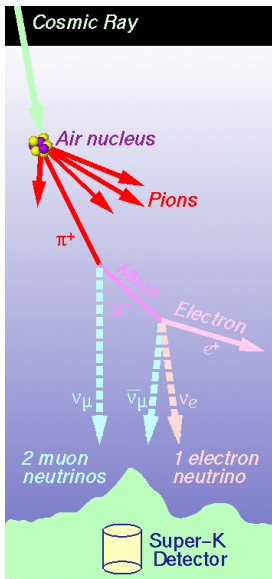
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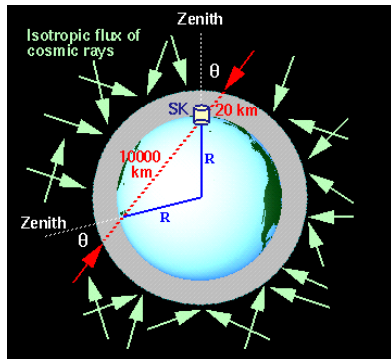
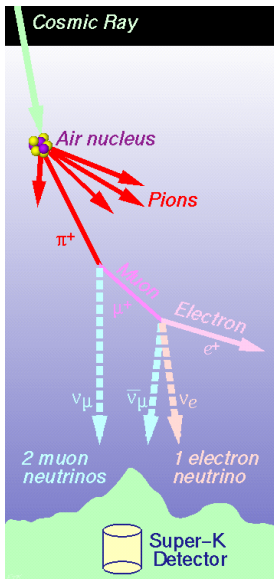
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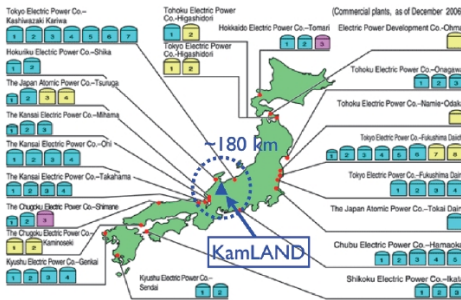


$L = 0$ to 13,000 km

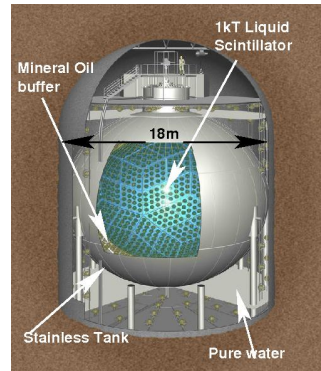
KamLAND: Reactor $\bar{\nu}_e \rightarrow \bar{\nu}_e$ Oscillations

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$$P(\bar{\nu}_e \rightarrow \bar{\nu}_x) = \sin^2 2\theta_{\text{solar}} \sin^2 \frac{1.27 \Delta m_{\text{solar}}^2 L}{E}$$



World reactors + Research reactors : 0.96%
Korean reactors : 3.2%



Two Different Mass Scales!

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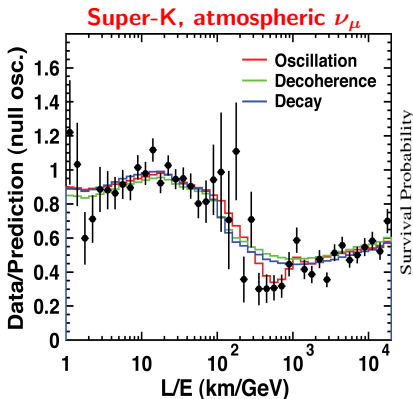
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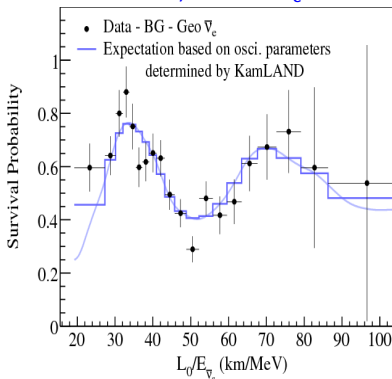
Global fit 2013:

$$\Delta m_{\text{atm}}^2 = 2.43_{-0.10}^{+0.06} \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{\text{atm}} = 0.386_{-0.21}^{+0.24}$$

Atmospheric L/E \sim 500 km/GeV

KamLAND, reactor $\bar{\nu}_e$



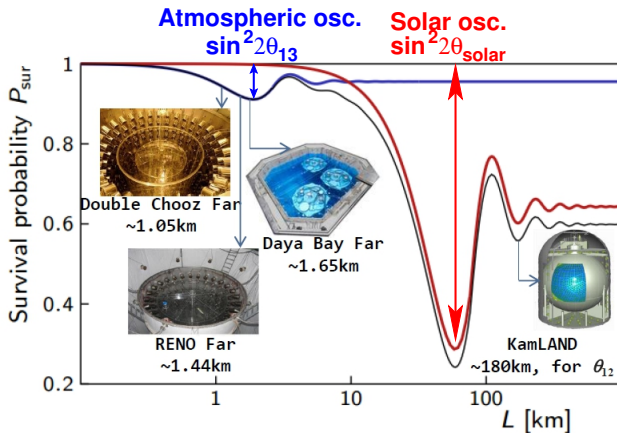
Global fit 2013:

$$\Delta m_{\text{solar}}^2 = 7.54_{-0.22}^{+0.26} \times 10^{-5} \text{ eV}^2$$

$$\sin^2 \theta_{\text{solar}} = 0.307_{-0.16}^{+0.18}$$

Solar L/E \sim 15,000 km/GeV

More Reactor $\bar{\nu}_e$: The 3rd Mixing Amplitude



In 2012: Measurement of a 3rd mixing amplitude:

$$\sin^2 \theta_{13} = 0.0241 \pm 0.0025$$

Much smaller than $\sin^2 \theta_{\text{solar}}$ and $\sin^2 \theta_{\text{atm.}}$

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Neutrino Mixing: 3 flavors, 3 amplitudes, 2 mass scales, and a CP violating phase

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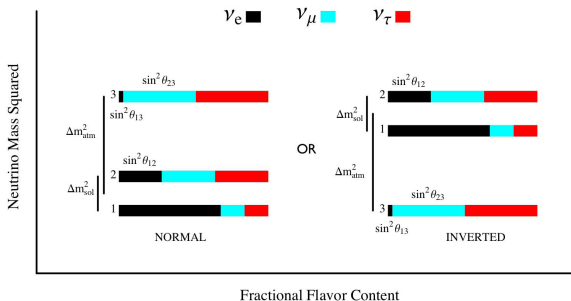
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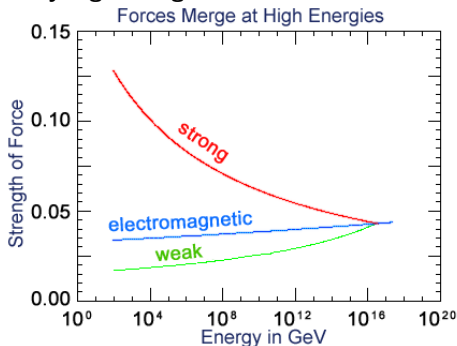


Parameter	Value (neutrino PMNS matrix)	Value (quark CKM matrix)
θ_{12} (solar)	$34 \pm 1^\circ$	$13.04 \pm 0.05^\circ$
θ_{23} (atm.)	$38 \pm 1^\circ$	$2.38 \pm 0.06^\circ$
θ_{13}	$8.9 \pm 0.5^\circ$	$0.201 \pm 0.011^\circ$
$\Delta m_{\text{solar}}^2$	$+(7.54 \pm 0.22) \times 10^{-5} \text{ eV}^2$	
$ \Delta m_{\text{atm.}}^2 $	$(2.43^{+0.10}_{-0.06}) \times 10^{-3} \text{ eV}^2$	
δ_{CP}	$-170 \pm 54^\circ$	$m_3 \gg m_2$ $67 \pm 5^\circ$

Unknown: Is $m_1 < m_3$ or vice versa?. Whats the value of δ_{cp} ?

Neutrino Masses in the Standard Model

In Grand Unified Theories all 3 forces (strong, electro-magnetic, weak) unify at very high energies:



A minimal re-normalizable extension to the Standard Model to account for 3 small left-handed neutrinos introduces **3 additional massive scalar right-handed sterile neutrinos** of mass $\sim M$. $M \sim \text{GUT scale or } M \sim 1 \text{ TeV (DARK MATTER?)}$

See-saw model: Neutrinos are Majorana ($\nu \equiv \bar{\nu}$)

Searching for CP Violation in Neutrino Oscillations

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Summary and
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Matter/anti-matter asymmetries in neutrinos are best probed using $\nu_\mu/\bar{\nu}_\mu \rightarrow \nu_e/\bar{\nu}_e$ oscillations (or vice versa):.

The charge-parity (CP) asymmetry is defined as

$$\mathcal{A}_{\text{cp}} = \frac{P(\nu_\mu \rightarrow \nu_e) - P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)}{P(\nu_\mu \rightarrow \nu_e) + P(\bar{\nu}_\mu \rightarrow \bar{\nu}_e)}$$

$$\mathcal{A}_{\text{cp}} \sim \frac{\cos \theta_{23} \sin 2\theta_{12} \sin \delta}{\sin \theta_{23} \sin \theta_{13}} \left(\frac{\Delta m_{21}^2 L}{4E_\nu} \right) + \text{matter effects}$$

from Z. Parsa, BNL

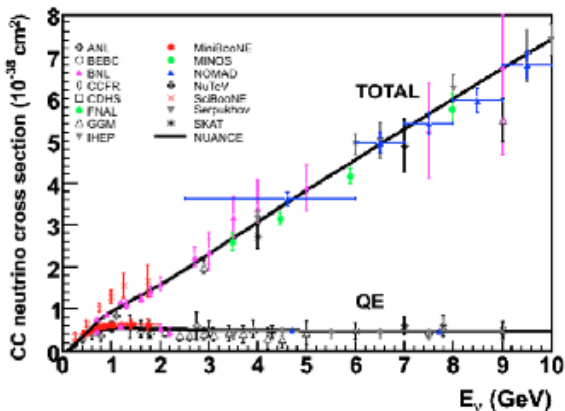
The CP phase δ_{cp} is unknown. CP is violated when $\delta_{\text{cp}} \neq 0, \pi$

The 3 most important things to know about ν CPV

- $\mathcal{A}_{\text{cp}} \propto 1/\sin \theta_{13} \Rightarrow$ Large θ_{13} makes CPV searches HARDER.
- $\mathcal{A}_{\text{cp}} \propto 1/E_\nu \Rightarrow$ CP asymmetries are larger at lower energies
- $\mathcal{A}_{\text{cp}} \propto L \Rightarrow$ CP asymmetries are larger at longer baselines

Neutrino Interaction Cross-sections

Neutrino CC cross-sections are very small and scale with energy:



Long baseline oscillations over 100's km and ≥ 200 MeV
energies needed to probe CPV

Matter Effect on Neutrino Oscillation

The Why,
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National
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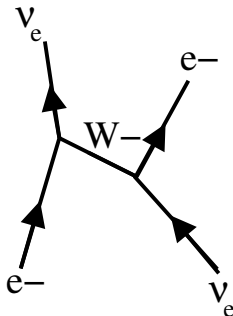
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1978 and 1986: L. Wolfenstein, S. Mikheyev and A. Smirnov propose the scattering of ν_e on electrons in matter adds a coherent forward scattering amplitude to neutrino oscillation amplitudes. This acts as a refractive index \Rightarrow **neutrinos in matter have different effective mass than in vacuum.**



The matter effect on ν_e scattering can be used to detect the unknown neutrino mass ordering using $\nu_x \rightarrow \nu_e$ oscillations through matter

Matter Effect on $P(\nu_\mu \rightarrow \nu_e)$

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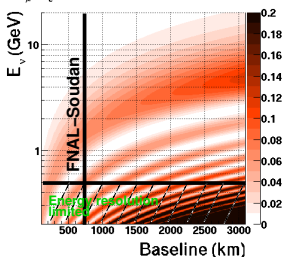
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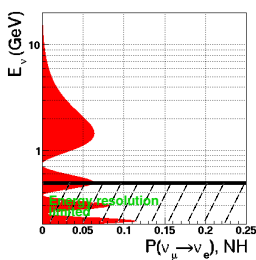
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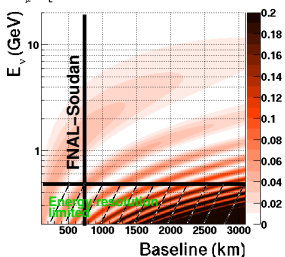
$P(\nu_\mu \rightarrow \nu_e), \text{NH}$



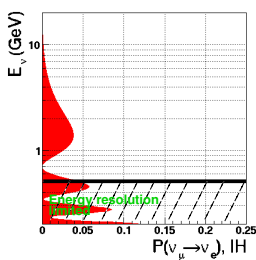
At 735km



$P(\nu_\mu \rightarrow \nu_e), \text{IH}$



At 735km



Matter Effect on $P(\nu_\mu \rightarrow \nu_e)$

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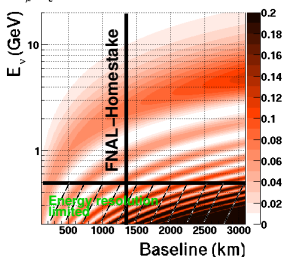
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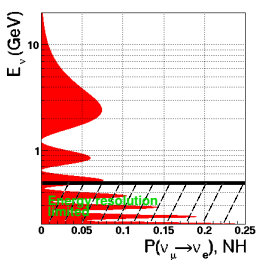
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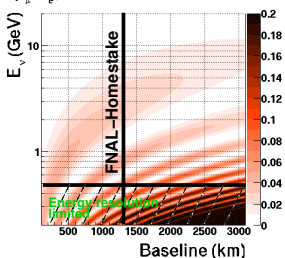
$P(\nu_\mu \rightarrow \nu_e), \text{NH}$



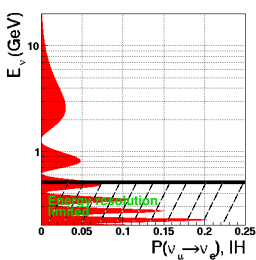
At 1300km



$P(\nu_\mu \rightarrow \nu_e), \text{IH}$



At 1300km



Matter Effect on $P(\nu_\mu \rightarrow \nu_e)$

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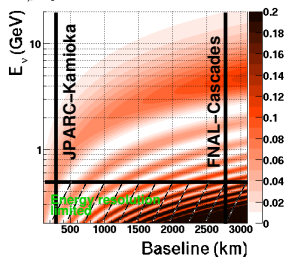
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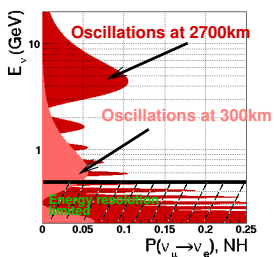
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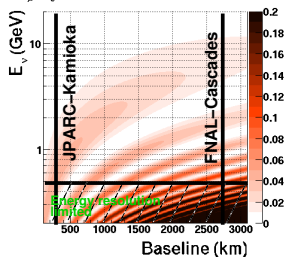
$P(\nu_\mu \rightarrow \nu_e), \text{NH}$



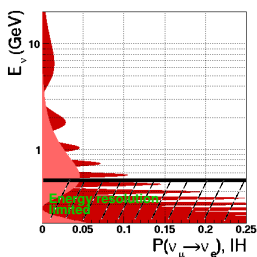
At 2700km



$P(\nu_\mu \rightarrow \nu_e), \text{IH}$

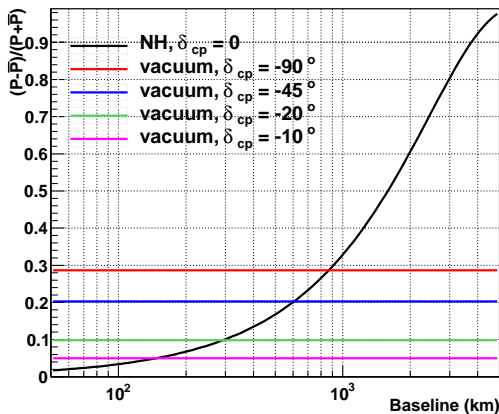


At 2700km



CP and Matter Asymmetries vs. Baselines for fixed L/E

CP asymmetries in $\nu_\mu \rightarrow \nu_e$ at 1st osc. node



Impact of the mass hierarchy on asymmetry is very large in long baseline experiments

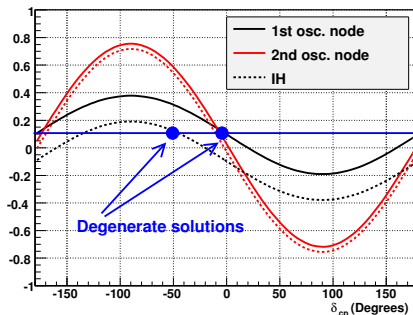
CP Asymmetries and the Mass Hierarchy

$\nu_\mu \rightarrow \nu_e$ oscillation maxima occur at

$$E_\nu^n (\text{GeV}) = \frac{2.5 \Delta m_{32}^2 (\text{eV}^2) L (\text{km})}{(2n - 1)\pi} \quad n = 1, 2, 3 \dots$$

$$L = 290 \text{ km}$$

Total Asymmetry at 290km



At short baselines, irreducible degeneracies with MH, δ_{cp}

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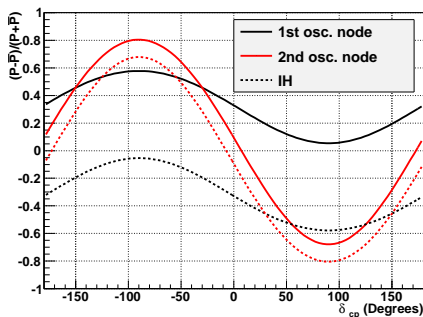
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$$E_\nu^n (\text{GeV}) = \frac{2.5 \Delta m_{32}^2 (\text{eV}^2) L (\text{km})}{(2n - 1)\pi} \quad n = 1, 2, 3 \dots$$

$$L = 1000 \text{ km}$$

Total Asymmetry at 1000km



A baseline $> 1000 \text{ km}$ is needed separate MH from δ_{cp}

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ACCELERATOR ν Experiments

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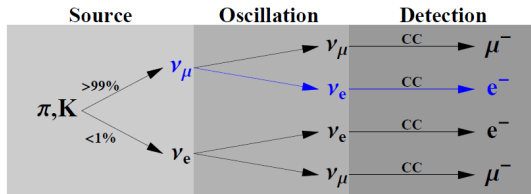
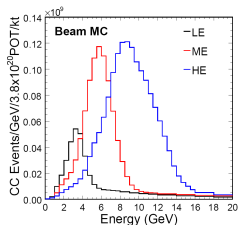
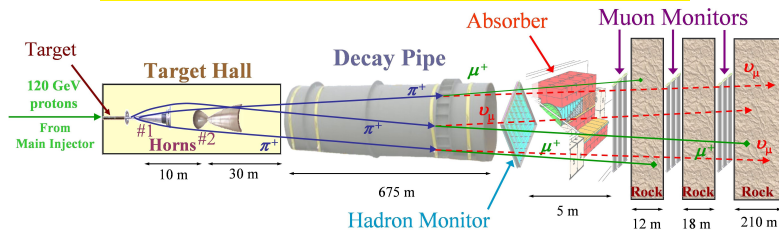
Summary and
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The Fermilab Accelerator Complex 2012



Neutrinos at the Main Injector (NuMI)

High power TUNABLE conventional neutrino beam:



NuMI as built cost (2004): ~ \$115 M (did not include all labor costs!!!)

> 3× more expensive in FY12 \$!

Fermilab Proton Plan

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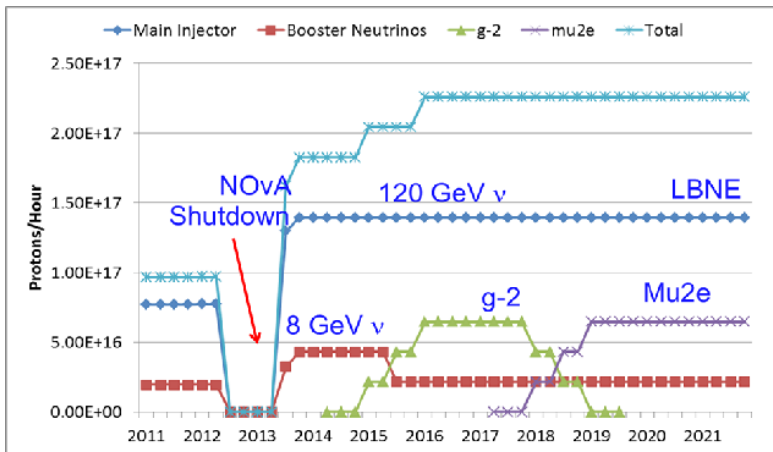
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Fermilab proton improvement plan: MI: 700 kW at 120 GeV by 2014

Future plans at Fermilab: Project X

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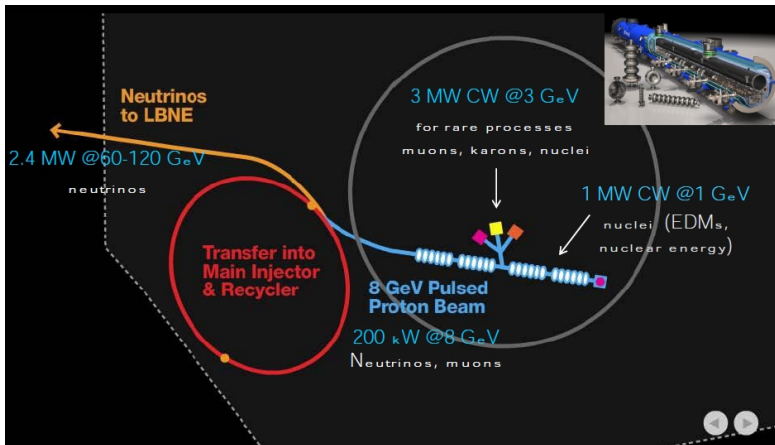
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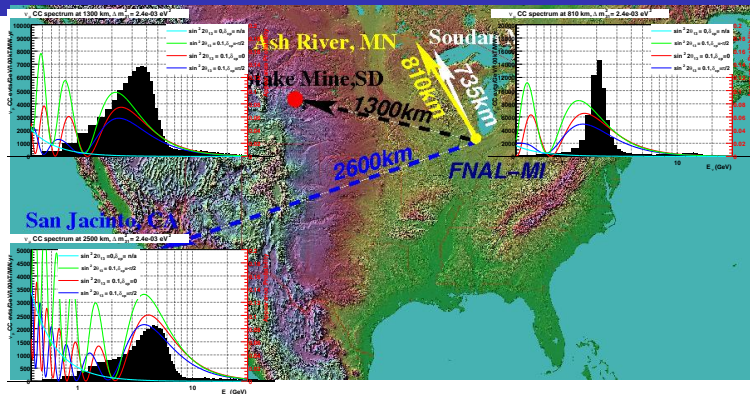
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Beyond 2025?: 2.4 MW at 60-120 GeV

Superbeam Baselines in the U.S.



CC event rates per 100kt.MW.yrs (1 MW.yr = 1×10^{21} p.o.t) for $\sin^2 2\theta_{13} = 0.1, \delta_{cp} = 0, \text{NH}$:

Expt	ν_μ CC	ν_μ CC osc	ν_μ NC	ν_e beam	$\nu_\mu \rightarrow \nu_e$	$\nu_\mu \rightarrow \nu_\tau$
Ash River 810km	18K	7.3K	3.6K	330	710	TBD
Hmstk 1300km	29K	11K	5.0K	280	1100	TBD
CA 2500km	11K	2.9K	1.6K	85	760	TBD

Need MW beams and 100 kton detectors regardless of baseline!

Massive Neutrino Detectors: Water Cherenkov

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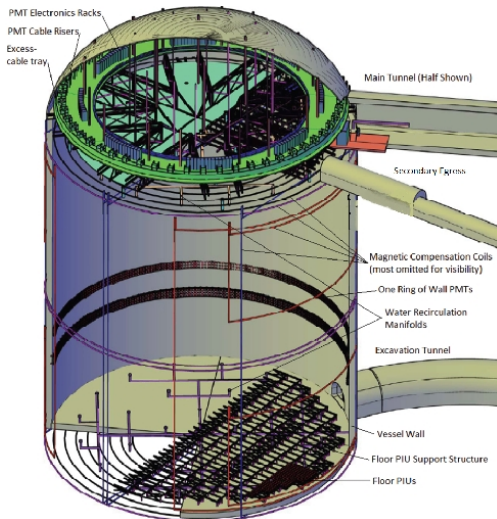
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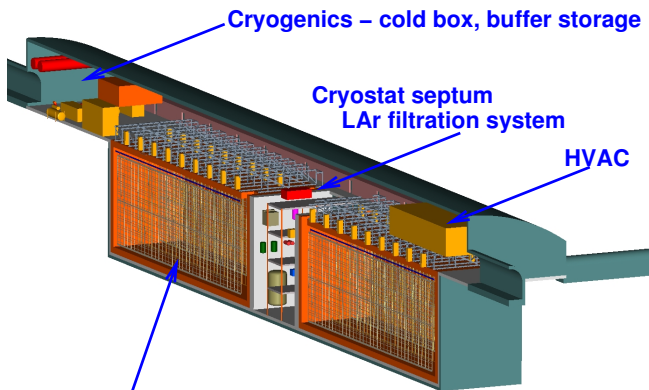
Summary and
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200kTon (underground) : ~ 700\$M FY10\$ (incl. 40% contingency)

Massive Neutrino Detectors: The LAr TPC

A 35 kton fiducial Liquid Argon Time-Projection-Chamber:



Detector Module
2 high x 3 wide x 18 long drift cells x 2 modules
216 APAs, 224 CPAs

35 kton (underground) : ~ 660\$M FY10\$ (incl. 40% contingency)

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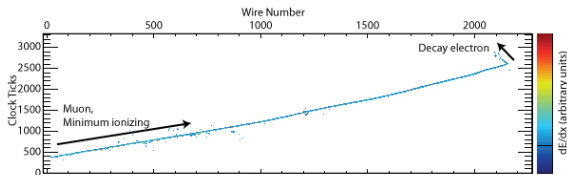
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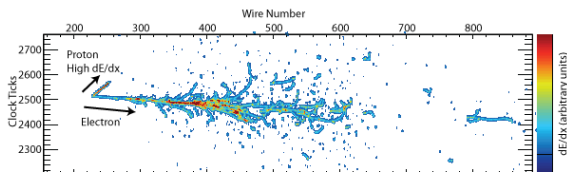
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Neutrino Interactions in a LAr-TPC

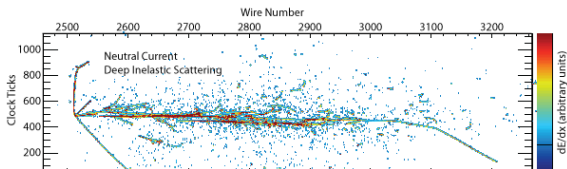
ν_μ CC



ν_e CC



ν_x NC



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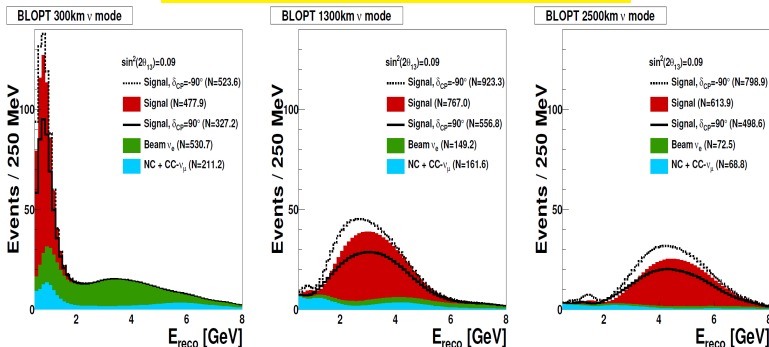
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Starting from the MI 120 GeV, we produced an optimized horn focused beam based on NuMI designs for each baseline. For shorter baselines we used off-axis angles. **For 35 kTon LAr-TPC.**

$\nu_\mu \rightarrow \nu_e$ Appearance. Normal Hierarchy



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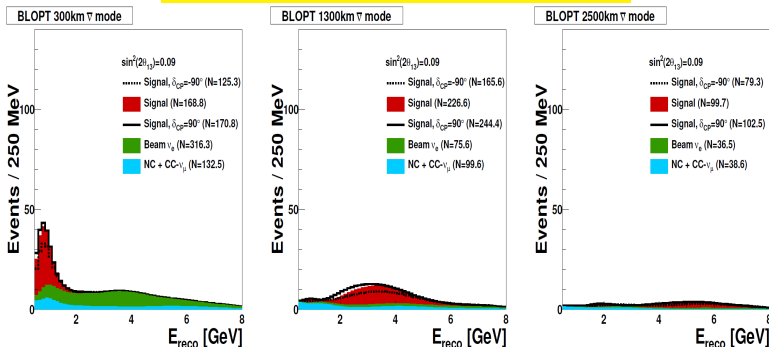
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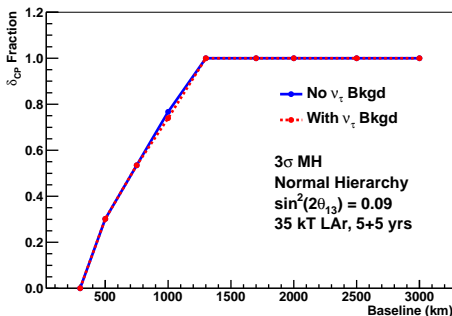
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Mass hierarchy sensitivity at $\geq 3\sigma$



Baselines of 1500-1700km sufficient to resolve MH

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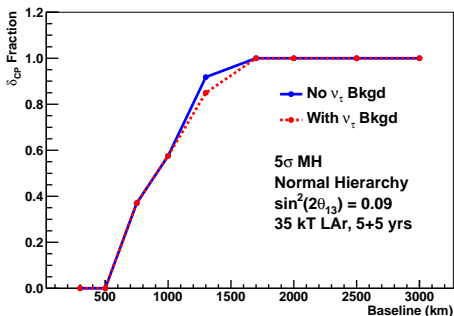
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Mass hierarchy sensitivity at $\geq 5\sigma$



Baselines of 1500-1700km sufficient to resolve MH

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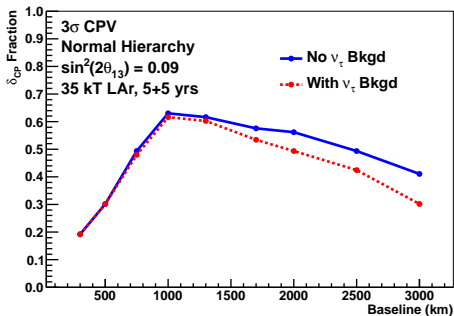
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Starting from the MI 120 GeV, we produced an optimized horn focused beam based on NuMI designs for each baseline. For shorter baselines we used off-axis angles. **For 35 kTon LAr-TPC.**

CP violation sensitivity at 3σ



Baselines of 1000-1500km best CPV sensitivity

The Long Baseline Neutrino Experiment

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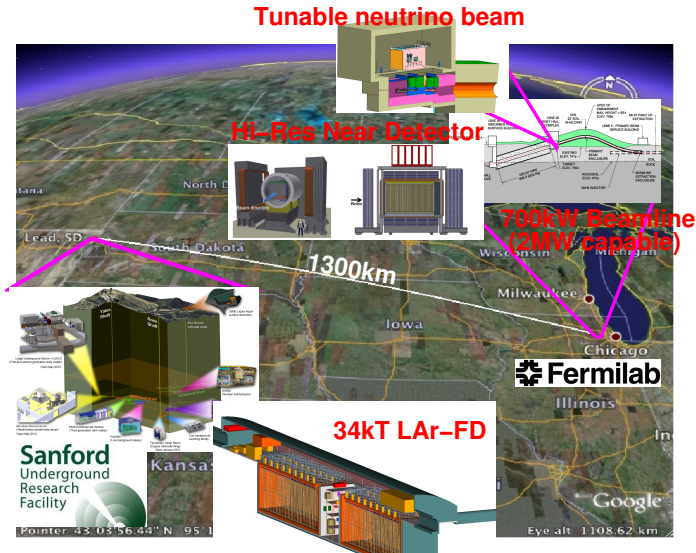
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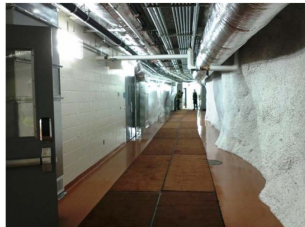


~ 350 people, 60 institutions from US, India, Italy, Japan, UK

The Sanford Underground Research Facility (SURF)

The Why,
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Dedicated underground research lab at Homestake Mine since 2008

Located 1300km from Fermilab : optimal baseline for CPV.

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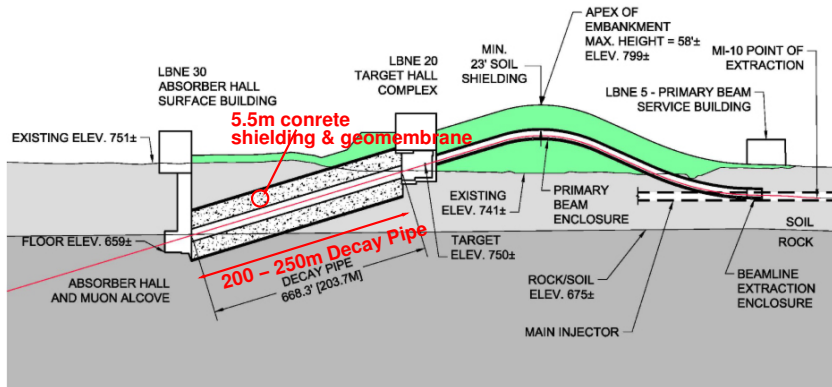
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The LBNE Beamline

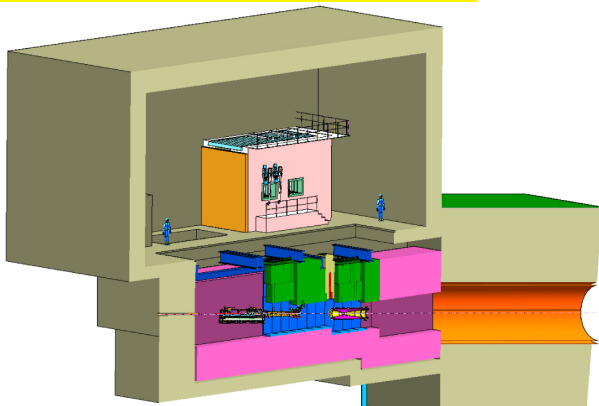
Novel beam-on-a-hill construction for 2.3MW from 60-120GeV



**Cost: ~ 390\$M AY (CD1) with 204m decay pipe
(incl 30% cont. and conventional facilities)**

The LBNE Beamline

Utilizes tunable NuMI targetry/focusing designs



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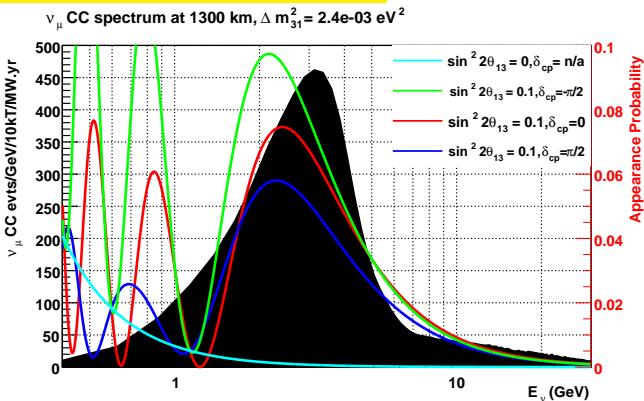
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The LBNE Beamline

Excellent match to $\nu_\mu \rightarrow \nu_e$ oscillation



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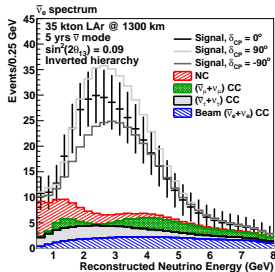
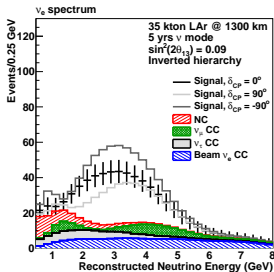
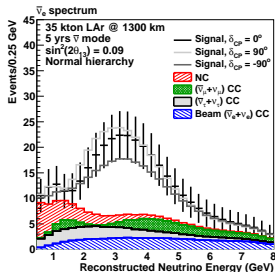
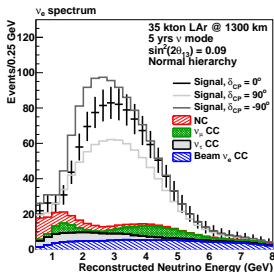
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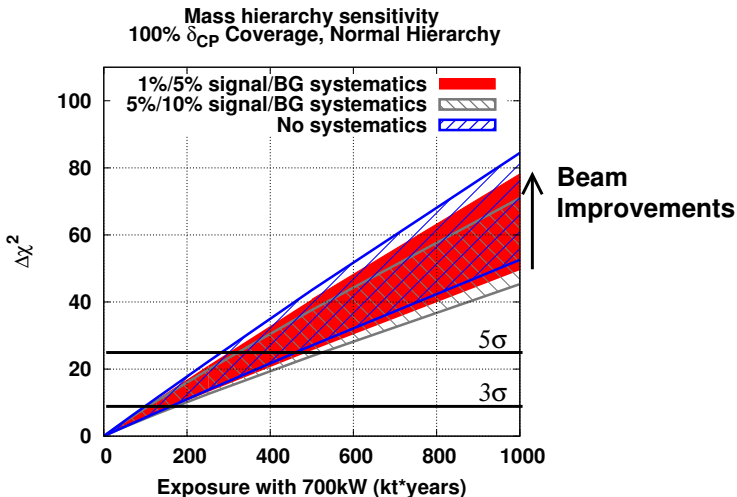
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LBNE MH/CPV Sensitivities with 700kW Beam, 35 kton.

M. Bass, CSU

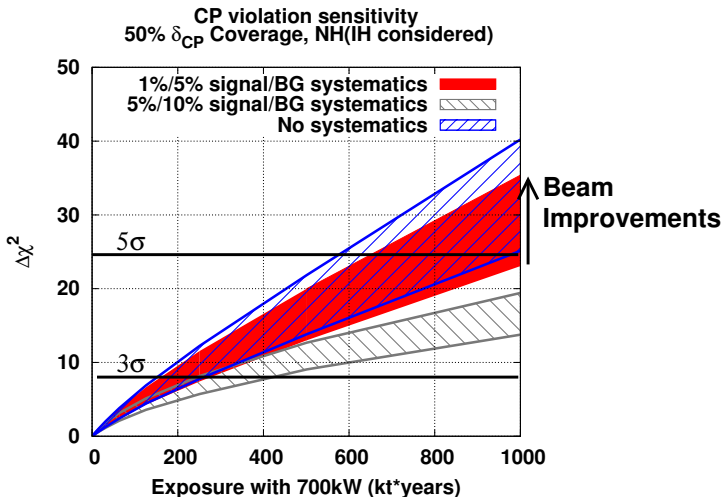
Worst case sensitivity to mass ordering vs exposure



LBNE MH/CPV Sensitivities with 700kW Beam, 35 kton.

M. Bass, CSU

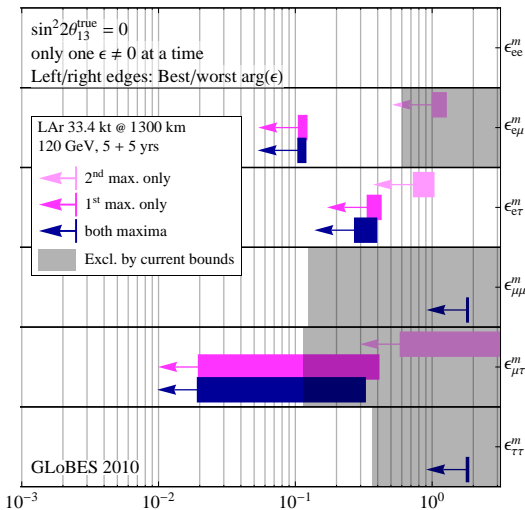
Worst case CPV sensitivity to 50% of δ_{CP} allowed values:



Sensitivity to New Physics, 35kTon, 700kW, 10yrs

Non-standard interactions (J. Kopp):

NC NSI discovery reach (3σ C.L.)



The Why,
What, Where
of the Long
Baseline
Neutrino
Experiment

Mary Bishai
Brookhaven
National
Laboratory

Why
Neutrinos:
History of ν s
Neutrino Mixing
Long Baseline ν
Oscillations

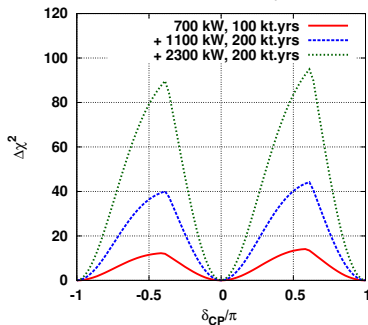
Accelerator ν
Experiments
Beams
Detectors

Which
Baseline?

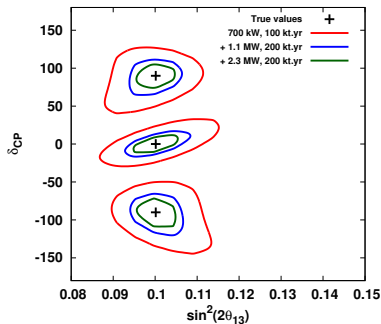
Whats LBNE?

Summary and
Conclusions

CPV significance vs δ_{CP}
Normal Hierarchy

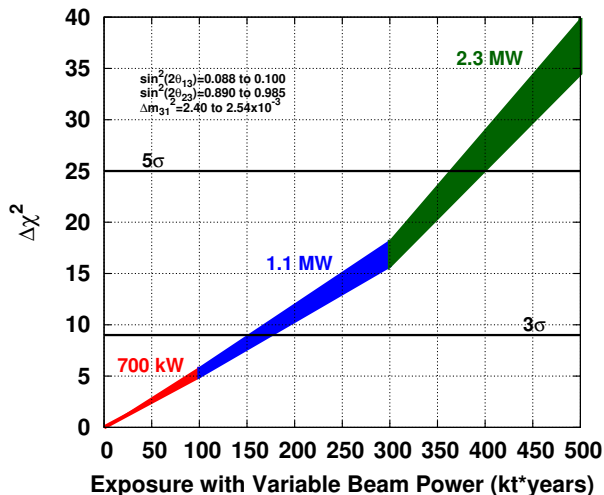


Project X Staging
1:1 $\nu:\bar{\nu}$, 1%/5% Signal/BG systematics



With sufficient exposure we can reach 5 – 10° resolution on δ_{CP} .

CP violation sensitivity 50% δ_{CP} Coverage, Normal Hierarchy



LBNE Phase I CD1 Approved Project

LBNE recieved DOE CD1 approval for phase I Dec 2012



Directed towards a distant detector at the Sanford Underground Research Facility (SURF) in Lead, SD

10 kton Liquid Argon TPC Far Detector just below the surface

All the Conventional Facilities at Fermilab and SURF required to support the beam and detectors.

DOE Goal: Phase I includes ND and > 10 kton underground by attracting non-DOE contributions.

LBNE CD1 Costs

J. Strait

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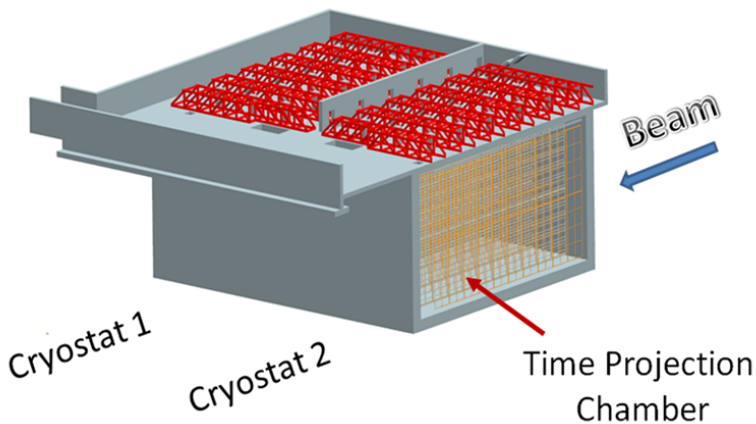
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LBNE Cost Estimate in At-Year k\$	Base Cost k\$	Contingency		TPC k\$	cost to date thru 6/2012
		k\$	% of cost to go		
130.01 Proj Office	57,014	16,099	32%	73,113	6,987
130.02 Beamline	129,370	35,311	29%	164,680	7,422
130.03 NDC	11,815	10,748	148%	22,563	4,558
130.04 WCD	11,178		0%	11,178	11,178
130.05 LAr FD	181,347	71,767	41%	253,114	7,775
130.06 CF	241,191	71,600	31%	312,790	6,872
Top-down Contingency		30,000	5%	30,000	
Grand Total	631,914	235,525	40%	867,439	44,792

Phase I: 10 kton LAr-TPC at SURF

10 kT Liquid Argon TPC for the LBNE Experiment:



Cost: ~ 260 \$ M AY (incl 40% cont, no CF), Ready 2022

Extra cost to go underground: ~ 135 M\$

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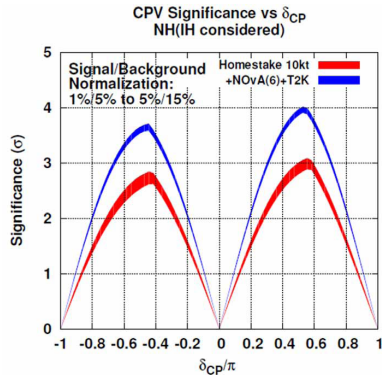
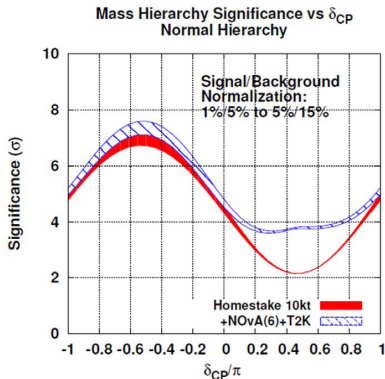
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The mass hierarchy can be determined using LBNE I+ T2K/NO ν A

LBNE I + T2K/NO ν A $\Rightarrow \geq 3\sigma$ sensitivity to CPV for 35% of δ_{CP}

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SUMMARY AND CONCLUSIONS

Summary

The 3 flavor framework for neutrino oscillations is well established

- Experiments have identified 2 very different mass scales and 3 mixing amplitudes.
- The mass ordering (hierarchy) of the 1 and 3 mass states is still unknown.
- The measurement of the hierarchy using the matter effect in long baseline accelerator $\nu_\mu \rightarrow \nu_e$ is the most effective technique.
MH determination in LB ν is independent of detector systematics
- Discovering ν CPV requires many 100kt .MW .yr exposures!
- While it is NOT necessary to know MH to measure CP violation . Baselines of > 1000 km are needed to cleanly separate CP violating effects from matter effects.
- The LBNE experiment has the 1) optimal baseline, 2) multi-MW beam design and 3) best detector technology. LBNE phase I preserves all 3 features. Extra non-DOE funding to increase phase I detector size, go underground, and add a near detector are being sought.

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Thank you